Gamma-Ray Bursts

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Atelier OCEVU « The physics of relativistic outflows », 22-24 March, 2016, IRAP, Toulouse

Emission from relativistic outflows: the case of Gamma-Ray Bursts Frédéric Daigne (Institut d'Astrophysique de Paris) Emission from relativistic outflows: the case of Gamma-Ray Bursts

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#### Introduction

- 1. GRBs: observed emission
- 2. Relativistic outflows in GRBs
- 3. Possible emission sites in GRBs
- 4. Modelling the emission from relativistic outflows

#### Some recent results

- 1. Weak photospheric emission in GRBs: constraints on magnetization
- 2. The origin of X-ray flares in GRB afterglows
- 3. Prompt GRB emission from internal shocks

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#### GRBs: observed emission





Prompt emission keV  $\rightarrow$  GeV (Fermi)

X-ray afterglow (Swift)



optical, GeV long-lasting Fermi/LAT emission

Observed prompt  $\gamma$ -ray spectrum

#### Fermi/GBM:

BB looked for in bright cases & found in many cases Fermi/LAT: 1st catalog extra-component in 4/28



#### Relativistic outflows in GRBs

Indirect: necessary to avoid a strong  $\gamma\gamma$  annihilation

Direct (in a few cases): apparent super-luminal motion

How relativistic are GRB outflows?

Pre-Fermi (MeV range) :  $\Gamma_{\rm min} \sim$  100-300

GeV detection by Fermi: stricter Lorentz factor constraints

- GRB 080916C:  $\Gamma_{\min} \ge 887$  (Abdo et al. 09)
- GRB 090510:  $\Gamma_{min} \ge 1200$  (Ackerman et al. 10)



# How relativistic are GRB outflows?

Detailed calculation: space/time/direction-dependent radiation field the estimate of  $\Gamma_{min}$  is reduced by a factor ~ 2-3 (see Granot et al. 2008; Hascoët, FD, Mochkovitch & Vennin 2012)



Model of bins a+b in GRB 080916C :  $\Gamma_{min} \sim 360$  (Hascoët et al. 2012) instead of ~900 (Abdo et al. 2009)

GRB 090926A: observed cutoff? See F. Piron's talk

## Apparent super-luminal motion in GRBs (radio afterglow)

#### Method 1:

Radio scintillation quenches as the source increases Transition diffractive / refractive : estimate of the angular size

#### Method 2:

0.5

0

(mas)

Declination

Relativ

GRB 030329

VLBI allows to resolve the late afterglow for nearby GRBs

From the size, the apparent velocity is deduced: superluminal apparent motion: relativistic motion

Apr6

Jun20

-0.5



Days After Burst

Taylor et al. 2004

0.5

0

Relative RA (mas)

# GRBs: possible emission sites





























## Modelling the emission from relativistic outflows

Modelling the emission from relativistic jets involves many steps:

- Dynamics
- Microphysics: magnetic field? Particle acceleration? Etc.
- Emission in the comoving frame?
- Emission in the observer frame (time delays, relativistic Doppler boosting, etc.)

Emission can be produced in several sites:

- Photosphere
- Internal shocks
- Reconnection
- External shock

Relativistic jets: dynamics (1)

In the most general case, dynamics of relativistic jets is complicated... (3D relativistic MHD ?)

In many cases, a simplified solution is possible :

- MHD? Not in the case of a non-magnetized outflow or a passive field
- Fluid? If the kinetic energy is dominant, a ballistic approach is possible (interacting shells, e.g. internal shock model for GRBs)

#### 3D/complex geometry?

- If the Lorentz factor is very large, global geometry is not important
  - $(1/\Gamma \ll \text{ jet opening angle})$
- If the Lorentz factor is very large, lateral expansion is negligible (c\_s/ $\Gamma \ll$  c)

Interaction with ambient medium/deceleration?

May be neglected when computing emission with an internal origin emitted well below the deceleration radius Relativistic jets: dynamics (2)

Usually, the initial conditions are not well known, especially due to the poor understanding of the physics of the relativistic ejection.

- Mass flux ?
- Energy flux ?
- Lorentz factor ?
- Magnetization / field geometry ?

# Relativistic jets: microphysics

It remains difficult to couple a dynamical calculation with a realistic microphysics. In addition, despite some recent progress, the relevant microphysics is not well understood (shock acceleration / reconnection).

This part is usually highly parametrized (e.g. "equipartition parameters", etc.)

# Relativistic jets: emission in the comoving frame

In the most general case, emission in the comoving frame is also complicated... (time-dependent radiative transfer with non-thermal particles)

- Many process: synchrotron, IC, γγ, etc.
- Leptons: primary electrons + secondary pairs due to  $\gamma\gamma$
- Contributions from hadrons?
- Two regimes:
  - radiatively efficient: radiative timescale  $\ll$  dynamical timescale
    - = particle cool immediately where they are accelerated, transport in the jet is not important
  - radiatively inefficient: radiative timescale  $\gg$  dynamical timescale
- Optical depth? Optically thin vs optically thick (comptonization ?)
   calculation is more complicated in the second case (multi-scatterings)
- Geometry of the photon field (important for IC, γγ, etc.) ?
   Etc.

## Relativistic jets: emission in the observer frame

- Integration over equal-arrival time surface (curvature of the emitting surface ?)
- Doppler effect
- (cosmological effects / interaction with ambient photon field, etc.)
- Complexity: multi-zones interactions
- e.g. photons emitted in a zone are scattered by accelerated electrons in another zone

# Weak quasi-thermal photospheric emission: constraints on the magnetization

e.g. GRB 120323A (short GRB) Guiriec [FD] et al. 2013



Warning: spectral analysis based on forward folding technique

e.g. GRB 080916C (long GRB) Guiriec [FD] et al. 2015



*Warning: spectral analysis based on forward folding technique* 

Non dissipative photosphere in magnetized outflows:

- Initial geometry is not specified
  Beyond R<sub>sph</sub>, the flow is radial (opening angle θ)
- Total injected power in the flow:  $\dot{E}$ -fraction  $\varepsilon_{th}$  is thermal -fraction 1- $\varepsilon_{th}$  is magnetic
- •Acceleration is complete at  $R_{sat} > R_{sph}$ •The final magnetization (above  $R_{sat}$ ) is  $\sigma$
- Photospheric emission occurs at R<sub>ph</sub>
   Non-thermal emission occurs above R<sub>ph</sub> with efficiency f<sub>NT</sub>

#### Three main parameters: $\epsilon_{th}$ , $\sigma$ , $f_{NT}$

Inversion method described by Pe'er et al. 2007 R<sub>0</sub>, R<sub>ph</sub>,  $\Gamma = F(data ; \epsilon_{th}, \sigma, f_{NT} ; z)$ 



**Fig. 1.** Schematic view of the problem geometry. The flow emerges from the central engine through a "circular opening" of radius  $\ell$ . Beyond a radius  $R_{sph}$  it expands radially within a cone of half opening  $\theta$ . The acceleration is completed at  $R_{sat}$ . The photosphere is located at  $R_{ph}$  and dissipation of kinetic and/or magnetic energy takes place at  $R_{diss}$ .

Hascoet, Daigne & Mochkovitch 2013

Three main parameters:  $\boldsymbol{\epsilon}_{\text{th}},\,\sigma,\,f_{\text{NT}}$ 

Inversion method described by Pe'er et al. 2007

$$R_0, R_{ph}, \Gamma = F(data ; \epsilon_{th}, \sigma, f_{NT} ; z)$$

$$\begin{split} R_0 &\simeq \left[ \frac{D_{\rm L} \mathcal{R}}{2(1+z)^2} \left( \frac{\phi}{1-\phi} \right)^{3/2} \right] \times \left[ \frac{f_{\rm NT}}{\epsilon_{\rm T}} \right]^{3/2}, \\ \Gamma &\simeq \left[ \frac{\sigma_{\rm T}}{m_{\rm p} c^3} \frac{(1+z)^2 D_{\rm L} F_{\rm BB}}{\phi} \frac{1-\phi}{\phi} \right]^{\frac{1}{4}} \times \left[ (1+\sigma) f_{\rm NT} \right]^{-1/4}, \\ R_{\rm ph} &\simeq \left[ \frac{\sigma_{\rm T}}{16m_{\rm p} c^3} \frac{D_{\rm L}^5 F_{\rm BB} \mathcal{R}^3}{(1+z)^6} \frac{1-\phi}{\phi} \right]^{\frac{1}{4}} \times \left[ (1+\sigma) f_{\rm NT} \right]^{-1/4}, \\ \phi &= F_{\rm BB}/F_{\rm tot} \qquad \mathcal{R} = \left( \frac{F_{\rm BB}}{\sigma T_{\rm BB}^4} \right)^{1/2}. \end{split}$$

Different scenarios:

- -Thermal acceleration (standard fireball):  $\epsilon_{th}$ =1 &  $\sigma$ =0 and f<sub>NT</sub> <10% (internal shocks)
- -Magnetized outflows:  $\varepsilon_{th} < 1$ 
  - -efficient acceleration:  $\sigma$  < 0.1-1 and f<sub>NT</sub> <10% (internal shocks)
  - -mag. outflow at large distance:  $\sigma >$  1 and  $\rm f_{NT} > 30\%$  (reconnection)

#### Exemple: GRB 100724B

Thermal component is weak (4% of total)



Guiriec [FD] et al. (2011)



Observations taken from Guiriec et al. (2011)

Exemple: GRB 100724B

Incompatible with the standard fireball, except for a very low  $R_0$  + very high non-thermal efficiency





Exemple: GRB 100724B

- Efficient magnetic acceleration + internal shocks
- Magnetized outflow at large distance + reconnection



#### Other exemples and summary

Most GRBs have a weak photosphere and are not compatible with the standard fireball :  $\epsilon_{\rm th} < 1\%$  (Daigne & Mochkovitch 2002)

•Exemples: GRB100724B (long) -non compatible with a standard fireball -compatible with efficient mag. acceleration + internal shocks ( $\epsilon_{th} < 1-10\%$ ) or magnetized outflow + reconnection (but low efficiency or  $\epsilon_{th} > 30\%$ ) (Guiriec et al. 2011; Hascoet et al. 2013)

GRB120323A (short) : similar conclusions, but allowing a larger  $\epsilon_{th}{>}50\%$ 

GRB 090902B: only case compatible with standard fireball

It implies a large initial magnetization in GRB outflows:

What is the magnetization  $\sigma$  at large distance? Internal dissipation by shocks or reconnection? X-ray flares produced by a long-lived RS: consequences for the structure of the relativistic ejecta



Afterglow from a long-lived RS: Genet, Daigne & Mochkovitch 2007 ; Uhm & Beloborodov 2007

# Afterglow from a long-lived RS

-long lived reverse shock: constraint on the initial Lorentz factor in the ejecta (Rees & Meszaros 98 ; Sari & Meszaros 00 ; Genet [FD] & Mochkovitch 07 ; Uhm & Beloborodov 07)

-dominant RS emission: constraint on microphysics RS vs FS ( $\epsilon_e$ ,  $\epsilon_B$ ) (Genet [FD] et al. 07; Uhm & Beloborodov 07; Uhm [FD] et al. 11)



# Afterglow from a long-lived RS

-long lived reverse shock: constraint on the initial Lorentz factor in the ejecta -dominant RS emission: constraint on microphysics RS vs FS ( $\epsilon_e$ ,  $\epsilon_B$ )

-No need for late energy injection to reproduce plateaus -Large diversity of lightcurves is expected (internal structure of the ejecta) (Uhm [FD] et al. 2012)

-Observed correlations between prompt and plateau properties can be reproduced (Hascoët [FD] et al. 2013)

-No need for late activity of the central engine to reproduce flares (Hascoët [FD] et al. arXiv:1503.08333)

## Long-lived RS afterglow: diversity

Top: FS (very low sensitivity to the internal structure of the ejecta)



Relativistic ejecta: constant  $10^{53}$  erg/s for 10 s - Source at z = 1FS:  $\epsilon_e = 10^{-2}$ ;  $\epsilon_B = 10^{-4}$ ; p = 2.3; RS:  $\epsilon_e = 10^{-1}$ ;  $\epsilon_B = 10^{-2}$ ; p = 2.3

Uhm [FD] et al. 2011

Each case corresponds to a different initial distribution of the Lorentz factor

-propagation of the reverse shock in a structured outflow -a signature of internal shocks?



An exemple of the distribution of Lorentz factor in the ejecta: (relativistic hydro simulation)

**---** Initial

---- During IS phase

End of IS phase (before deceleration)

Hascoët et al. astro-ph/1503.08333

-propagation of the reverse shock in a structured outflow -a signature of internal shocks?



Hascoët et al. astro-ph/1503.08333

-propagation of the reverse shock in a structured outflow -a signature of internal shocks?



Hascoët et al. astro-ph/1503.08333

-propagation of the reverse shock in a structured outflow -a signature of internal shocks?



Flares are produced when the RS crosses a dense shell formed in the IS phase Hascoët et al. astro-ph/1503.08333

-propagation of the reverse shock in a structured outflow -a signature of internal shocks?



Hascoët et al. astro-ph/1503.08333

Needs hydro simulation + radiative calculation for validation

-propagation of the reverse shock in a structured outflow

-a signature of internal shocks?



internal shocks, what is the contribution of these shocks to the prompt emission? Sub-photospheric dissipation or direct emission?

#### Prompt gamma-ray emission from internal shocks?

How to distinguish between the proposed mechanisms for the prompt emission?

- -Lighcurves: OK for all scenarios
- -Spectrum
- -Spectral evolution





Low-energy photon index in fast cooling synchrotron spectrum?

-3/2 : pure fast cooling synchrotron ~ -1 : fast cooling synchrotron + inverse Compton in KN regime (Derishev et al. 01 ; Bosnjak et al. 09 ; Wang et al. 09 ; Daigne et al. 11) -2/3 : marginally fast cooling synchrotron (Daigne et al. 11 ; Beniamini & Piran 13) -1 → -0.5 : fast cooling synchrotron + IC in decaying magnetic field (Derishev 07 ; Lemoine 13 ; Uhm & Zhang 14 ; Zhao et al. 14)



- Band vs Band+BB: different low-energy photon index? Compatible with (modified) fast cooling synchrotron?
   e.g. GRB120323A α=-0.92 → -1.4 Guiriec [FD] et al. 2013 GRB 080916C α=-1.0 → -1.2 Guiriec [FD] et al. 2015 etc.
- Inconsistency between time-integrated and time-resolved analysis?
- Shape of the extra-component in LAT is not well constrained. Is the X-ray excess real?







Band function used both in time-integrated/resolved analysis





e.g. New analysis of GRB090926A with Pass 8 (LAT photons  $\times$  2.4) : see F. Piron's talk



Ackermann et al. 2011: Band (steep  $\alpha$ ) + PL (with cutoff in bin c) – X-ray excess





Yassine, Piron, Daigne & Mochkovitch, Fermi Symposium 2015





Example of a simulated GRB pulse produced by internal shocks (full simulation: dynamics+radiation)



Light curve in BATSE range : channels 1 (blue) to 4 (red)

Example of a simulated GRB pulse produced by internal shocks (full simulation: dynamics+radiation)



Extra component

Bosnjak & Daigne 2014

Example of a simulated GRB pulse produced by internal shocks (full simulation: dynamics+radiation)





Preece et a.l. 2014

Not shown: hardness-intensity correlation slope 1.4

#### Prompt GeV emission from internal shocks



Bosnjak & Daigne 2014 ; see also Asano & Meszaros.

# Summary

# Summary

In the scenario where the prompt GRB emission is produced above the photosphere:

The weak quasi-thermal photospheric emission implies a high magnetization at the base of the relativistic outflow.

•Then, depending on the magnetization at large distance, internal dissipation responsible of the prompt emission can be either shocks or reconnection.

•When deceleration by the external medium starts, the reverse shock may have an important contribution to the afterglow emission (constraints on Lorentz factor + microphysics). It can explain the complexity/diversity of the afterglow light-curves without strong assumptions for the central engine lifetime/energetics.

In this scenario, X-ray flares can be a signature of previous internal shocks.

•Then internal shocks may be responsible for the prompt emission, under some strong constraints on the microphysics.